Topological Quantum Field Theory for Calabi-Yau threefolds and G_2 -manifolds

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1. Introduction

In the past two decades we witness many fruitful interactions between mathematics and physics. One example is the Donaldson-Floer theory for oriented four manifolds. Physical considerations leads to the discovery of the Seiberg-Witten theory which has profound impact to our understandings of four manifolds. Another example is the Mirror Symmetry for Calabi-Yau manifolds. This duality transformation in the string theory leads to many surprising predictions in the enumerative geometry.

String theory in physics studies a ten dimensional space-time $X \times \mathbb{R}^{3,1}$ with X a six dimensional Riemannian manifold with its holonomy group inside SU (3), the so-called Calabi-Yau threefold. Certain parts of the Mirror Symmetry conjecture, as studied by Vafa's group, are specific for Calabi-Yau manifolds of complex dimension three. They include the Gopakumar-Vafa conjecture on Gromov-Witten invariants of arbitrary genus for Calabi-Yau threefolds, the Ooguri-Vafa conjecture on the relationships between knot invariants and enumerations of holomorphic disks and so on. The key reason is they belong to dualities for G_2 -manifolds. G_2 -manifolds can be naturally interpreted as special Octonion manifolds [23]. For any Calabi-Yau threefold X, the seven dimensional manifold $X \times S^1$ is automatically a G_2 -manifold because of the natural inclusion SU (3) $\subset G_2$.

In recent years, there are many studies of G_2 -manifolds in M-theory including works of Archaya, Atiyah, Gukov, Vafa, Witten, Yau, Zaslow and many others (e.g. [1], [5], [13], [2]).

In the studies of the symplectic geometry of a Calabi-Yau threefold X, we consider unitary flat bundles over three dimensional (special) Lagrangian submanifolds L in X. The corresponding theory for a G_2 -manifold M is called the *special* \mathbb{H} -Lagrangian geometry (or C-geometry in [19]). where we consider Anti-Self-Dual (abbrev. ASD) bundles over four dimensional coassociative submanifolds, or equivalently special \mathbb{H} -Lagrangian submanifolds of type II [23], (abbrev. \mathbb{H} -SLag) C in M.

Counting ASD bundles over a fixed four manifold C is the well-known theory of Donaldson differentiable invariants, Don(C). Similarly, counting unitary flat bundles over a fixed three manifold L is Floer's Chern-Simons homology theory $HF_{CS}(L)$. When C is a connected sum $C_1\#_LC_2$ along a homology three sphere, the relative Donaldson invariants $Don(C_i)$'s take values in $HF_{CS}(L)$ and

 $Don\left(C\right)$ can be recovered from individual pieces by a gluing theorem, $Don\left(C\right) = \langle Don\left(C_1\right), Don\left(C_2\right)\rangle_{HF_{CS}(L)}$ (see e.g. [7]). Similarly when L has a handlebody decomposition $L = L_1 \#_{\Sigma} L_2$, each L_i determines a Lagrangian subspace \mathcal{L}_i in the moduli space $\mathcal{M}^{flat}\left(\Sigma\right)$ of unitary flat bundles over the Riemann surface Σ and Atiyah conjectures that we can recover $HF_{CS}\left(L\right)$ from the Floer's Lagrangian intersection homology group of \mathcal{L}_1 and \mathcal{L}_2 in $\mathcal{M}^{flat}\left(\Sigma\right)$, $HF_{CS}\left(L\right) = HF_{Lag}^{\mathcal{M}^{flat}\left(\Sigma\right)}\left(\mathcal{L}_1,\mathcal{L}_2\right)$. Such algebraic structures in the Donaldson-Floer theory can be formulated as a Topological Quantum Field Theory (abbrev. TQFT), as defined by Segal and Atiyah [3].

In this paper we propose a TQFT by counting ASD bundles over four dimensional \mathbb{H} -SLags C in any closed (almost) G_2 -manifold M, called \mathbb{H} -SLags cycles. They can be identified as zeros of a naturally defined closed one form on the configuration space of topological cycles. We expect to obtain a homology theory $H_C(M)$ by applying the Witten's Morse theory. When M is non-compact with an asymptotically cylindrical end $X \times [0,\infty)$, then the set of boundary data of relative \mathbb{H} -SLag cycles determines a Lagrangian submanifold \mathcal{L}_M in the moduli space $\mathcal{M}^{SLag}(X)$ of special Lagrangian cycles in the Calabi-Yau threefold X.

When we decompose $M = M_1 \#_X M_2$ along an infinite asymptotically cylindrical neck, it is reasonable to expect to have a gluing formula,

$$H_{C}\left(M\right)=HF_{Lag}^{\mathcal{M}^{SLag}\left(X\right)}\left(\mathcal{L}_{M_{1}},\mathcal{L}_{M_{2}}\right).$$

The main technical difficulty in defining this TQFT rigorously is the *compactness* issue for the moduli space of \mathbb{H} -SLags in M. We do not know how to resolve this problem and our homology groups are only defined in the *formal* sense (and physical sense?).

2. G_2 -manifolds and \mathbb{H} -SLag geometry

We first review some basic definitions and properties of G_2 -geometry, see [19] for more details.

Definition 1. A seven dimensional Riemannian manifold M is called a G_2 -manifold if the holonomy group of its Levi-Civita connection is inside $G_2 \subset SO(7)$.

The simple Lie group G_2 can be identified as the subgroup of SO(7) consisting of isomorphism $g: \mathbb{R}^7 \to \mathbb{R}^7$ preserving the linear three form Ω ,

$$\Omega = f^{1} f^{2} f^{3} - f^{1} \left(e^{1} e^{0} + e^{2} e^{3} \right) - f^{2} \left(e^{2} e^{0} + e^{3} e^{1} \right) - f^{3} \left(e^{3} e^{0} + e^{1} e^{2} \right),$$

where $e^0, e^1, e^2, e^3, f^1, f^2, f^3$ is any given orthonormal frame of \mathbb{R}^7 . Such a three form, or up to conjugation by elements in $GL(7, \mathbb{R})$, is called *positive*, and it determines a unique compatible inner product on \mathbb{R}^7 [6].

Gray [12] shows that G_2 -holonomy of M can be characterized by the existence of a positive harmonic three form Ω .

Definition 2. A seven dimensional manifold M equipped with a positive closed three form Ω is called an almost G_2 -manifold.

Remark: The relationship between G_2 -manifolds and almost G_2 -manifolds is completely analogous to the relationship between Kahler manifolds and symplectic manifolds.

For example, suppose that X is a complex three dimensional Kähler manifold with a trivial canonical line bundle, i.e. there exists a nonvanishing holomorphic

three form Ω_X . Yau's celebrated theorem says that there is a Kähler form ω_X on X with holonomy in SU(3), i.e. a Calabi-Yau threefold. In particular both Ω_X and ω_X are parallel forms. Then the product $M=X\times S^1$ is a G_2 -manifold with

$$\Omega = \operatorname{Re} \Omega_X + \omega_X \wedge d\theta.$$

Conversely, one can prove, using Bochner arguments, every G_2 -metric on $X \times S^1$ must be of this form. More generally, if ω_X is a general Kähler form on X, then $(X \times S^1, \Omega)$ is an almost G_2 -manifold and the converse is also true.

Next we quickly review the geometry of \mathbb{H} -SLag cycles in an almost G_2 -manifold (see [19]).

DEFINITION 3. An orientable four dimensional submanifold C in an almost G_2 -manifold (M,Ω) is called a coassociative submanifold, or simply a \mathbb{H} -SLag, if the restriction of Ω to C is identically zero,

$$\Omega|_C=0.$$

If M is a G_2 -manifold, then such a C is calibrated by $*\Omega$ in the sense of Harvey and Lawson [14], in particular, it is an absolute minimal submanifold in M. The normal bundle of any \mathbb{H} -SLag C can be naturally identified with the bundle of self-dual two forms on C. McLean [27] shows that infinitesimal deformations of any \mathbb{H} -SLag are unobstructed and they are parametrized by the space of harmonic self-dual two forms on C, i.e. $H^2_+(C,\mathbb{R})$.

For example, if S is a complex surface in a Calabi-Yau threefold X, then $S \times \{t\}$ is a \mathbb{H} -SLag in $M = X \times S^1$ for any $t \in S^1$. Similarly, if L is a three dimensional special Lagrangian submanifold in X with phase $\pi/2$, i.e. $\omega|_L = \operatorname{Re} \Omega_X|_L = 0$, then $L \times S^1$ is also a \mathbb{H} -SLag in $M = X \times S^1$.

Definition 4. A \mathbb{H} -SLag cycle in an almost G_2 -manifold (M,Ω) is a pair (C,D_E) with C a \mathbb{H} -SLag in M and D_E an ASD connection over C.

Remark: H-SLag cycles are supersymmetric cycles in physics as studied in [26]. Their moduli space admits a natural three form and a cubic tensor [19], which play the roles of correlation function and Yukawa coupling in physics.

We assume that the ASD connection D_E over C has rank one, i.e. a U(1) connection. This avoids the occurrence of reducible connections, thus $\mathcal{M}^{\mathbb{H}-SLag}(M)$ is a smooth manifold. It has a natural orientation and its expected dimension equals $b^1(C)$, the first Betti number of C. This is because the moduli space of \mathbb{H} -SLags has dimension equals $b_+^2(C)$ [27] and the existence of an ASD U(1)-connection over C is equivalent to $H_-^2(C, \mathbb{R}) \cap H^2(C, \mathbb{Z}) \neq \phi$. The number $b^1(C)$ is responsible for twisting by a flat U(1)-connection.

For simplicity, we assume that $b^1(C) = 0$, otherwise, one can cut down the dimension of $\mathcal{M}^{\mathbb{H}-SLag}(M)$ to zero by requiring the ASD connections over C to have trivial holonomy around loops $\gamma_1, ..., \gamma_{b^1(C)}$ in C representing an integral basis of $H_1(C, \mathbb{Z})$. We plan to count the algebraic number of points in this moduli space $\#\mathcal{M}^{\mathbb{H}-SLag}(M)$.

This number, in the case of $X \times S^1$, can be identified with a proposed invariant of Joyce [17] defined by counting rigid special Lagrangian submanifolds in any Calabi-Yau threefold. To explain this, we need the following proposition on the strong rigidity of product \mathbb{H} -SLags.

PROPOSITION 5. If $L \times S^1$ is a \mathbb{H} -SLag in $M = X \times S^1$ with X a Calabi-Yau threefold, then any \mathbb{H} -SLag representing the same homology class must also be a product.

Proof: For simplicity we assume that the volume of the S^1 factor is unity, $Vol\left(S^1\right)=1$. If $L\times S^1$ is a \mathbb{H} -SLag in M then L is special Lagrangian submanifold in X with phase $\pi/2$, i.e. $\operatorname{Re}\Omega_X|_L=\omega|_L=0$. Suppose C is another \mathbb{H} -SLag in M representing the same homology class, we have $Vol\left(C\right)=Vol\left(L\right)$. If we write $C_\theta=C\cap(X\times\{\theta\})$ for any $\theta\in S^1$, then $Vol\left(C_\theta\right)\geq Vol\left(L\right)$, as L is a calibrated submanifold in X. Furthermore the equality sign holds only if C_θ is also calibrated. In general we have

$$Vol(C) \ge \int_{S^1} Vol(C_{\theta}) d\theta,$$

with the equality sign holds if and only if C is a product with S^1 . Combining these, we have

$$Vol\left(L\right) = Vol\left(C\right) \geq \int_{S^{1}} Vol\left(C_{\theta}\right) d\theta \geq \int_{S^{1}} Vol\left(L\right) d\theta = Vol\left(L\right).$$

Thus both inequalities are indeed equal. Hence $C=L'\times S^1$ for some special Lagrangian submanifold L' in X.

Suppose $M=X\times S^1$ is a product G_2 -manifold and we consider product \mathbb{H} -SLag $C=L\times S^1$ in M. From the above proposition, every \mathbb{H} -SLag representing [C] must also be a product. Since $b_+^2(C)=b^1(L)$, the rigidity of the \mathbb{H} -SLag C in M is equivalent to the rigidity of the special Lagrangian submanifold L in X. When this happens, i.e. L is a rational homology three sphere, we have $b^2(C)=0$ and

No. of ASD U(1)-bdl/
$$C=\#H^2\left(C,\mathbb{Z}\right)=\#H^2\left(L,\mathbb{Z}\right)=\#H_1\left(L,\mathbb{Z}\right).$$

Here we have used the fact that the first cohomology group is always torsion free. Thus the number of such \mathbb{H} -SLag cycles in $X \times S^1$ equals the number of special Lagrangian rational homology three spheres in a Calabi-Yau threefold X, weighted by $\#H_1(L,\mathbb{Z})$. Joyce [17] shows that with this particular weight, the numbers of special Lagrangians in any Calabi-Yau threefold behave well under various surgeries on X, and expects them to be invariants. Thus in this case, we have

$$\#\mathcal{M}^{\mathbb{H}-SLag}(X\times S^1) = \text{Joyce's proposed invariant for } \#\text{SLag. in } X.$$

In the next section, we will propose a homology theory, whose Euler characteristic gives $\#\mathcal{M}^{\mathbb{H}-SLag}(M)$.

3. Witten's Morse theory for H-SLag cycles

We are going to use the parametrized version of \mathbb{H} -SLag cycles in any almost G_2 -manifold M. We fix an oriented smooth four dimensional manifold C and a rank r Hermitian vector bundle E over C. We consider the *configuration space*

$$C = Map(C, M) \times A(E).$$

DEFINITION 6. An element (f, D_E) in C is called a parametrized \mathbb{H} -SLag cycles in M if

$$f^*\Omega = F_E^+ = 0,$$

where the self-duality is defined using the pullback metric from M.

Remark: From the positivity of Ω , $f^*\Omega = 0$ implies f is an immersion.

Instead of Aut(E), the symmetry group \mathcal{G} in our situation consists of gauge transformations of E which cover arbitrary diffeomorphisms on M,

$$\begin{array}{ccc}
E & \stackrel{g}{\rightarrow} & E \\
\downarrow & & \downarrow \\
M & \stackrel{g_M}{\rightarrow} & M
\end{array}$$

It fits into the following exact sequence,

$$1 \to Aut(E) \to \mathcal{G} \to Diff(C) \to 1.$$

The natural action of \mathcal{G} on \mathcal{C} is given by

$$g \cdot (f, D_E) = (f \circ g_M, g^*D_E),$$

for any $(f, D_E) \in \mathcal{C} = Map(C, M) \times \mathcal{A}(E)$. Notice that \mathcal{G} preserves the set of parametrized \mathbb{H} -SLag cycles in M.

The configuration space \mathcal{C} has a natural one form Φ_0 : At any $(f, D_E) \in \mathcal{C}$ we can identify the tangent space of \mathcal{C} as

$$T_{(f,D_E)}\mathcal{C} = \Gamma\left(C, f^*T_M\right) \times \Omega^1\left(C, ad\left(E\right)\right).$$

We define

$$\Phi_{0}\left(f,D_{E}\right)\left(v,B\right)=\int_{C}Tr\left[f^{*}\left(\iota_{v}\Omega\right)\wedge F_{E}+f^{*}\Omega\wedge B\right],$$

for any $(v, B) \in T_{(f, D_E)} \mathcal{C}$.

PROPOSITION 7. The one form Φ_0 on C is closed and invariant under the action by G.

Proof: Recall that there is a universal connection \mathbb{D}_E over $C \times \mathcal{A}(E)$ whose curvature \mathbb{F}_E at a point (x, D_E) equals,

$$\mathbb{F}_{E}|_{(x,D_{E})} = \left(\mathbb{F}_{E}^{2,0}, \mathbb{F}_{E}^{1,1}, \mathbb{F}_{E}^{0,2}\right) \\
\in \Omega^{2}\left(C\right) \otimes \Omega^{0}\left(\mathcal{A}\right) + \Omega^{1}\left(C\right) \otimes \Omega^{1}\left(\mathcal{A}\right) + \Omega^{0}\left(C\right) \otimes \Omega^{2}\left(\mathcal{A}\right)$$

with

$$\mathbb{F}_{E}^{2,0} = F_{E}, \, \mathbb{F}_{E}^{1,1} \, (v,B) = B \, (v) \, , \, \mathbb{F}_{E}^{0,2} = 0,$$

where $v \in T_xC$ and $B \in \Omega^1(C, ad(E)) = T_{D_E}A(E)$ (see e.g. [20]). The Bianchi identity implies that $Tr\mathbb{F}_E$ is a closed form on $C \times A(E)$. We also consider the evaluation map,

$$ev: C \times Map(C, M) \rightarrow M$$

 $ev(x, f) = f(x)$.

It is not difficult to see that the pushforward of the differential form $ev^*(\Omega) \wedge Tr\mathbb{F}_E$ on $C \times Map(C, M) \times \mathcal{A}(E)$ to $Map(C, M) \times \mathcal{A}(E)$ equals Φ_0 , i.e.

$$\Phi_{0}=\int_{C}ev^{st}\left(\Omega
ight)\wedge Tr\mathbb{F}_{E}.$$

Therefore the closedness of Φ_0 follows from the closedness of Ω . It is also clear from this description of Φ_0 that it is \mathcal{G} -invariant.

From now on, we assume that E is a rank one bundle.

LEMMA 8. The zeros of Φ_0 are the same as parametrized \mathbb{H} -SLag cycles in M.

Proof: Suppose (f, D_E) is a zero of Φ_0 . By evaluating it on various (0, B), we have $f^*\Omega = 0$, i.e. $f: C \to M$ is a parametrized \mathbb{H} -SLag. This implies that the map

$$\Box \Omega: T_{f(x)}M \to \Lambda^2 T_x^*C$$

has image equals $\Lambda_+^2 T_x^* C$, for any $x \in C$. By evaluating Φ_0 on various (v,0), we have $F_E^+ = 0$, i.e. (f, D_E) is a parametrized \mathbb{H} -SLag cycle in M. The converse is obvious.

From above results, Φ_0 descends to a closed one form on \mathcal{C}/\mathcal{G} , called Φ . Locally we can write $\Phi = d\mathcal{F}$ for some function \mathcal{F} whose critical points are precisely (unparametrized) \mathbb{H} -SLag cycles in M. Using the gradient flow lines of \mathcal{F} , we could formally define a Witten's Morse homology group, as in the famous Floer's theory. Roughly speaking one defines a complex (\mathbf{C}_*, ∂) , where \mathbf{C}_* is the free Abelian group generated by critical points of \mathcal{F} and ∂ is defined by counting the number of gradient flow lines between two critical points of relative index one.

Remark: The equations for the gradient flow are given by

$$\frac{\partial f}{\partial t} = * (f^* \xi \wedge F_E), \frac{\partial D_E}{\partial t} = * (f^* \Omega),$$

where $\xi \in \Omega^{2}\left(M,T_{M}\right)$ is defined by $\left\langle \xi\left(u,v\right),w\right\rangle =\Omega\left(u,v,w\right)$. The equation

$$\partial^2 = 0$$

requires a good compactification of the moduli space of \mathbb{H} -SLag cycles in M, which we are lacking at this moment. We denote this proposed homology group as $H_C(M)$, or $H_C(M, \alpha)$ when $f_*[C] = \alpha \in H_4(M, \mathbb{Z})$.

This homology group should be invariant under deformations of the almost G_2 -metric on M and its Euler characteristic equals,

$$\chi\left(H_{C}\left(M
ight)
ight)=\#\mathcal{M}^{\mathbb{H}-SLag}\left(M
ight).$$

Like Floer homology groups, they measure the *middle dimensional* topology of the configuration space C divided by G.

4. TQFT of H-SLag cycles

In this section we study complete almost G_2 -manifold M_i with asymptotically cylindrical ends and the behavior of $H_C(M)$ when a closed almost G_2 -manifold M decomposes into connected sum of two pieces, each with an asymptotically cylindrical end,

$$M = M_1 \# M_2.$$

Nontrivial examples of compact G_2 -manifolds are constructed by Kovalev [18] using such connected sum approach. The boundary manifold X is necessary a Calabi-Yau threefold. We plan to discuss analytic aspects of M_i 's in a future paper [24].

Each M_i 's will define a Lagrangian subspace \mathcal{L}_{M_i} in the moduli space of special Lagrangian cycles in X. Furthermore we expect to have a gluing formula expressing the above homology group for M in terms of the Floer Lagrangian intersection homology group for the two Lagrangian subspaces \mathcal{L}_{M_1} and \mathcal{L}_{M_2} ,

$$H_{C}\left(M\right)=HF_{Lag}^{\mathcal{M}^{SLag}\left(X\right)}\left(\mathcal{L}_{M_{1}},\mathcal{L}_{M_{2}}\right).$$

These properties can be reformulated to give us a topological quantum field theory. To begin we have the following definition.

DEFINITION 9. An almost G_2 -manifold M is called cylindrical if $M = X \times \mathbb{R}^1$ and its positive three form respect such product structure, i.e.

$$\Omega_0 = \operatorname{Re} \Omega_X + \omega_X \wedge dt.$$

A complete almost G_2 -manifold M with one end $X \times [0, \infty)$ is called asymptotically cylindrical if the restriction of its positive three form equals to the above one for large t, up to a possible error of order $O(e^{-t})$. More precisely the positive three form Ω of M restricted to its end equals,

$$\Omega = \Omega_0 + d\zeta$$

 $\label{eq:constraint} \textit{for some two form } \zeta \; \textit{satisfying} \; |\zeta| + |\nabla \zeta| + \left| \nabla^2 \zeta \right| + \left| \nabla^3 \zeta \right| \leq C e^{-t}.$

Remark: If M is an almost G_2 -manifold with an asymptotically cylindrical end $X \times [0, \infty)$, then (X, ω_X, Ω_X) is a complex threefold with a trivial canonical line bundle, but the Kähler form ω_X might not be Einstein. This is so, i.e. a Calabi-Yau threefold, provided that M is a G_2 -manifold. We will simply write $\partial M = X$.

We consider $\mathbb{H}\text{-SLags }C$ in M which satisfy a Neumann condition at infinity. That is, away from some compact set in M, the immersion $f:C\to M$ can be written as

$$f:L\times [0,\infty)\to X\times [0,\infty)$$

with $\partial f/\partial t$ vanishes at infinite [24]. A relative \mathbb{H} -SLag itself has asymptotically cylindrical end $L \times [0, \infty)$ with L a special Lagrangian submanifold in X. A relative \mathbb{H} -SLag cycle in M is a pair (C, D_E) with C a relative \mathbb{H} -SLag in M and D_E a unitary connection over C with finite energy,

$$\int_C |F_E|^2 dv < \infty.$$

Any finite energy connection D_E on C induces a unitary flat connection $D_{E'}$ on L [7].

Such a pair $(L, D_{E'})$ of a unitary flat connection $D_{E'}$ over a special Lagrangian submanifold L in a Calabi-Yau threefold X is called a *special Lagrangian cycle* in X. Their moduli space $\mathcal{M}^{SLag}(X)$ plays an important role in the Strominger-Yau-Zaslow Mirror Conjecture [28] or [22]. The tangent space to $\mathcal{M}^{SLag}(X)$ is naturally identified with $H^2(L,\mathbb{R})\times H^1(L,ad(E'))$. For line bundles over L, the cup product

$$\cup: H^2(L,\mathbb{R}) \times H^1(L,\mathbb{R}) \to \mathbb{R},$$

induces a symplectic structure on $\mathcal{M}^{SLag}\left(X\right)$ [15]. Using analytic results from [24] about asymptotically cylindrical manifolds, we can prove the following theorem.

CLAIM 10. Suppose M is an asymptotically cylindrical (almost) G_2 -manifold with $\partial M = X$. Let $\mathcal{M}^{\mathbb{H}-SLag}(M)$ be the moduli space of rank one relative \mathbb{H} -SLag cycles in M. Then the map defined by the boundary values,

$$b: \mathcal{M}^{\mathbb{H}-SLag}(M) \to \mathcal{M}^{SLag}(X)$$
,

is a Lagrangian immersion.

Sketch of the proof ([24]): For any closed Calabi-Yau threefold X (resp. G_2 -manifold M), the moduli space of rank one special Lagrangian submanifolds L (resp. \mathbb{H} -SLags C) is smooth [27] and has dimension $b^2(L)$ (resp. $b_+^2(C)$). The same holds true for complete manifold M with a asymptotically cylindrical end $X \times [0, \infty)$, where $b_+^2(C)_{L^2}$ denote the dimension of L^2 -harmonic self-dual two forms on a relative \mathbb{H} -SLag C in M.

The linearization of the boundary value map $\mathcal{M}^{\mathbb{H}-SLag}(M) \to \mathcal{M}^{SLag}(X)$ is given by $H^2_+(C)_{L^2} \stackrel{\alpha}{\to} H^2(L)$. Similar for the connection part, where the boundary value map is given by $H^1(C)_{L^2} \stackrel{\beta}{\to} H^1(L)$. We consider the following diagram where (i) each row is a long exact sequence of L^2 -cohomology groups for the pair (C,L) and (ii) each column in a perfect pairing.

Notice that $H^2_+(C, L)$, $H^2_+(C)$ and $H^2(L)$ parametrize infinitesimal deformation of C with fixed ∂C , deformation of C alone and deformation of L respectively.

By simply homological algebra, it is not difficult to see that $\operatorname{Im}\alpha \oplus \operatorname{Im}\beta$ is a Lagrangian subspace of $H^{2}\left(L\right) \oplus H^{1}\left(L\right)$ with the canonical symplectic structure. Hence the result. \blacksquare

We denote the immersed Lagrangian submanifold $b\left(\mathcal{M}^{\mathbb{H}-SLag}\left(M\right)\right)$ in $\mathcal{M}^{SLag}\left(X\right)$ by \mathcal{L}_{M} . When M decompose as a connected sum $M_{1}\#_{X}M_{2}$ along a long neck, as in Atiyah's conjecture on Floer Chern-Simons homology group [3], we expect to have an isomorphism,

$$H_C(M) \cong HF_{Lag}^{\mathcal{M}^{SLag}(X)}(\mathcal{L}_{M_1}, \mathcal{L}_{M_2}).$$

More precisely, suppose Ω_t with $t \in [0, \infty)$, is a family of G_2 -structure on $M_t = M$ such that as t goes to infinite, M decomposes into two components M_1 and

 M_2 , each has an aymptotically cylindrical end $X \times [0, \infty)$. Then we expect that $\lim_{t\to\infty} H_C(M_t) \cong HF_{Lag}^{\mathcal{M}^{SLag}(X)}(\mathcal{L}_{M_1}, \mathcal{L}_{M_2})$. We summarize these structures in the following table:

Manifold:	(almost) G_2 -manifold, M^7	(almost) CY threefold, X^6	
SUSY Cycles:	H-SLag. submfds.+ ASD bundles	SLag submfds.+ flat bundles	
Invariant:	Homology group, $H_{C}\left(M\right)$	Fukaya category, $Fuk\left(\mathcal{M}^{SLag}\left(X\right)\right)$.	

These associations can be formalized to form a TQFT [4]. Namely we associate an additive category $F(X) = Fuk\left(\mathcal{M}^{SLag}\left(X\right)\right)$ to a closed almost Calabi-Yau threefold X, a functor $F\left(M\right): F\left(X_{0}\right) \to F\left(X_{1}\right)$ to an almost G_{2} -manifold M with asymptotically cylindrical ends $X_{1}-X_{0}=X_{1}\cup\bar{X}_{0}$. They satisfy

- (i) $F(\phi)$ = the additive tensor category of vector spaces ((Vec)),
- (ii) $F(X_1 \coprod X_2) = F(X_1) \otimes F(X_2)$.

For example, when M is a closed G_2 -manifold, that is a cobordism between empty manifolds, then we have $F(M): ((Vec)) \to ((Vec))$ and the image of the trivial bundle is our homology group $H_C(M)$.

5. More TQFTs

There are other TQFTs naturally associated to Calabi-Yau threefolds and G_2 -manifolds but (1) they do not involve nontrivial coupling between submanifolds and bundles and (2) new difficulties arise because of corresponding moduli spaces for Calabi-Yau threefolds have virtual dimension zero and could be singular. They are essentially in the paper by Donaldson and Thomas [9].

TQFT of associative cycles

We assume that M is a G_2 -manifold, i.e. Ω is parallel rather than closed. Three dimensional submanifolds A in M calibrated by Ω is called associative submanifolds and they can be characterized by $\chi|_A=0$ ([14]) where $\chi\in\Omega^3(M,T_M)$ is defined by $\langle w,\chi(x,y,z)\rangle=*\Omega(w,x,y,z)$. We define a parametrized A-cycle to be a pair $(f,D_E)\in\mathcal{C}_A=Map(A,M)\times\mathcal{A}(E)$, with $f:A\to M$ a parametrized A-submanifold and D_E is a unitary flat connection on a Hermitian vector bundle E over A. There is also a natural \mathcal{G} -invariant closed one form Φ_A on \mathcal{C}_A given by

$$\Phi_{A}\left(f,D_{E}\right)\left(v,B\right)=\int_{A}TrF_{E}\wedge B+\left\langle f^{*}\chi,v\right\rangle _{T_{M}},$$

for any $(v,B) \in \Gamma(A,f^*T_M) \times \Omega^1(A,ad(E)) = T_{(f,D_E)}\mathcal{C}_A$. Its zero set is the moduli space of A-cycles in M. As before, we could formally apply arguments in Witten's Morse theory to Φ_A and define a homology group $H_A(M)$.

The corresponding category associated to a Calabi-Yau threefold X would be the Fukaya-Floer category of the moduli space of unitary flat bundles over holomorphic curves in X, denote $\mathcal{M}^{curve}(X)$. We summarize these in the following

table:

Manifold: G_2 -manifold, M^7		CY threefold, X^6	
SUSY Cycles:	A-submfds.+ flat bundles	Holomorphic curves+ flat bundles	
Invariant:	Homology group, $H_A(M)$	Fukaya category, $Fuk\left(\mathcal{M}^{curve}\left(X\right)\right)$.	

TQFT of Donaldson-Thomas bundles

We assume that M is a seven manifold with a G_2 -structure such that its positive three form Ω is co-closed, rather than closed, i.e. $d\Theta = 0$ with $\Theta = *\Omega$. In [9] Donaldson and Thomas introduce a first order Yang-Mills equation for G_2 -manifolds,

$$F_E \wedge \Theta = 0.$$

Their solutions are the zeros of the following gauge invariant one form Φ_{DT} on $\mathcal{A}(E)$,

$$\Phi_{DT}\left(D_{E}\right)\left(B\right) = \int_{M} Tr\left[F_{E} \wedge B\right] \wedge \Theta,$$

for any $B \in \Omega^1(M, ad(E)) = T_{D_E} \mathcal{A}(E)$. This one form Φ_{DT} is closed because of $d\Theta = 0$. As before, we can formally define a homology group $H_{DT}(M)$. The corresponding category associated to a Calabi-Yau threefold X should be the Fukaya-Floer category of the moduli space of Hermitian Yang-Mills connections over X, denote $\mathcal{M}^{curve}(X)$. Again we summarize these in a table:

Manifold:	G_2 -manifold, M^7	CY threefold, X^6	
SUSY Cycles:	DT-bundles	Hermitian YM-bundles	
Invariant:	Homology group, $H_{DT}(M)$	Fukaya category, $Fuk\left(\mathcal{M}^{HYM}\left(X\right)\right)$.	

It is an interesting problem to understand the transformations of these TQFTs under dualities in M-theory.

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